

Review:

Thm: (second order sufficient cond.)

An extremal $y(\cdot)$ is a strict weak minimum of the Basic VP if $L_{zz}(x, y, y') > 0$, $\forall x \in [a, b]$ and the interval $[a, b]$ contains no points conjugate to a .

Proof: y is extremal $\Rightarrow \delta J|_y(\eta) = 0$

$$\Rightarrow J(y + \alpha\eta) = J(y) + \delta^2 J|_y(\eta) \alpha^2 + o(\alpha^2)$$

and we already argued that $\delta^2 J|_y(\eta) > 0$ for $\eta \neq 0$.

So, we just need to argue that $\delta^2 J|_y(\eta)$ dominates $o(\alpha^2)$.

$\rightarrow$ Recall that no conjugate point to a on $[a, b]$ implies existence of a $W(x): [a, b] \rightarrow \mathbb{R}$ satisfying Riccati equation $P(Q + W') = W^2$ on $[a, b]$,

$$(I) \quad \text{s.t.} \quad \delta^2 J|_y[\eta] = \int_a^b \underbrace{P(x)}_{>0} \underbrace{\left(\eta'(x) + \frac{W(x)}{P(x)} \eta(x) \right)^2}_{>0 \text{ if } \eta \neq 0} dx.$$

Now, because $P(x) = L_{zz}(x) > 0$

$P(x) > 0$ on $[a, b] \rightarrow \exists \varepsilon > 0$ s.t. $P(x) > \varepsilon > 0$, $\forall x \in [a, b]$.

$$(I) \Rightarrow \int_a^b \left[(P(x) - \varepsilon) (\eta'(x))^2 + Q(x) \cdot \eta(x)^2 \right] dx > 0,$$

for all $\eta \neq 0$ admissible

$$\Rightarrow \int_a^b \left[P(x) (\eta'(x))^2 + Q(x) \cdot (\eta(x))^2 \right] dx > \varepsilon \int_a^b (\eta'(x))^2 dx$$

because (show this!) $\int_a^b \eta^2(x) dx \leq \frac{(b-a)^2}{2} \int_a^b (\eta'(x))^2 dx$

$$\Rightarrow \delta^2 J[\eta] = \int_a^b [P(x) \cdot (\eta'(x))^2 + Q(x) \cdot (\eta(x))^2] dx > c \cdot \varepsilon \left[\int_a^b ((\eta')^2 + \eta^2) dx \right]$$

for some positive constant $c = \frac{1}{1 + (b-a)^2/2}$.

Recall from (**) for $\frac{c}{2} \varepsilon > 0$, $\exists \delta > 0$ s.t.

$$|\bar{P}(x, y, y', \alpha)| < \frac{c}{2} \varepsilon, \quad |\bar{Q}(x, y, y', \alpha)| < \frac{c}{2} \varepsilon \quad \forall x, \quad \forall \eta, \quad \forall |\alpha| < \delta.$$

$$\Rightarrow \left| \frac{O(\alpha^2)}{\alpha^2} \right| = \left| \int_a^b [\bar{P}(\eta')^2 + \bar{Q} \eta^2] dx \right| \leq \frac{c}{2} \varepsilon \left[\int_a^b ((\eta')^2 + \eta^2) dx \right]$$

$$\Rightarrow \delta^2 J[\eta] \cdot \alpha^2 + O(\alpha^2) \geq \frac{c}{2} \varepsilon \cdot \left[\int_a^b ((\eta')^2 + \eta^2) dx \right] \quad \square$$

Note: • this argument will fail for α -norm $\|\eta\|_0$, so
one sufficient condition is only for weak minimum.

• For the vector case, everything is similar except that

$$P(x) = \frac{1}{2} L_{zz} > 0 \quad \text{P.D. on } [a, b]$$

positive definite matrix

and the Riccati-type equation reads

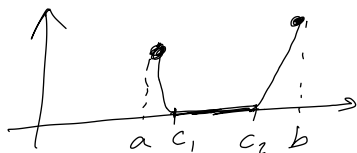
$$\boxed{Q + W' = W P^{-1} W}$$

Chap. 3: From cal. of variations to optimal control!

→ So far, we obtained conditions for weak extremums y to different variational problem: $y \in C^1([a,b]) \Rightarrow y \in C^1([a,b])$.

→ In optimal control (and other setting), we are more interested to strong solutions. Here, we focus on the piecewise C^1 curves: $y \in C^0([a,b])$ and y is C^1 except finitely many points

examples: 1) The catenary (chain) problem when the ground is too close to the hanging points compared to the length of the chain:



@ c_1 and c_2 the curve is not C^1 !

2) Imagine in controlling a car, sudden break to avoid collisions; a jump in acceleration implies a non-smoothness in the velocity curve!

Def: we call the points where y is not C^1 , corner points.

Q: Suppose y is a piecewise- C^1 curve and is a strong extremal for the Basic VP. what do we know about y ?

1) it is a weak extremal between each two corner points.
 $\Rightarrow$ must satisfy E-L eq.

more formally, it should be a weak^{*} solution w.r.t.

the generalized 1-norm:

$$\|y\|_1 \triangleq \max_{x \in [a, b]} |y(x)| + \max_{x \in [a, b]} \max \{ |y'(x^-)|, |y'(x^+)| \}.$$

2) we will obtain two new conditions for strong extremums:

I) weierstrass - Erdmann corner condition for strong extremum y :

$p = L_z$ and $H = p \cdot y' - L$ (momentum and Hamiltonian)
 are continuous at corner points.

II) we can also generalize the Legendre's condition to
corner points for a strong minimum:

$$L_{zz}(x, y(x), y^\pm(x)) \geq 0.$$

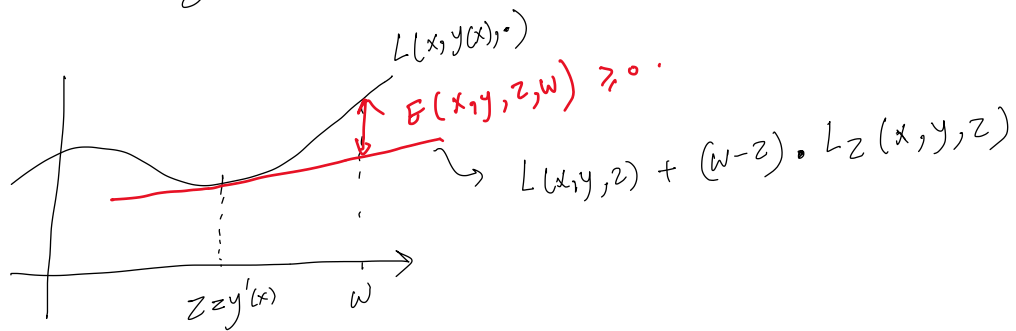
III) weierstrass - excess condition for strong minimum y :

$$\text{define } E(x, y, z, w) \triangleq L(x, y, w) - L(x, y, z) - (w - z) \cdot L_z(x, y, z)$$

for $x \in [a, b]$ non-corner points.

then $E(x, y(x), y'(x), w) \geq 0$, $\forall x \in [a, b]$
 $\forall w \in \mathbb{R}^n$

Geometrically, this means $L(x, y(x), \cdot)$ is locally convex at $y'(x)$.



Exercise: show that in Hamiltonian terminology

$$E(x, y, y', w) = H(x, y, y', p) - H(x, y, w, p)$$

where $p = L_z(x, y, y')$ is the momentum. And,
 by treating p and y' as independent variables,

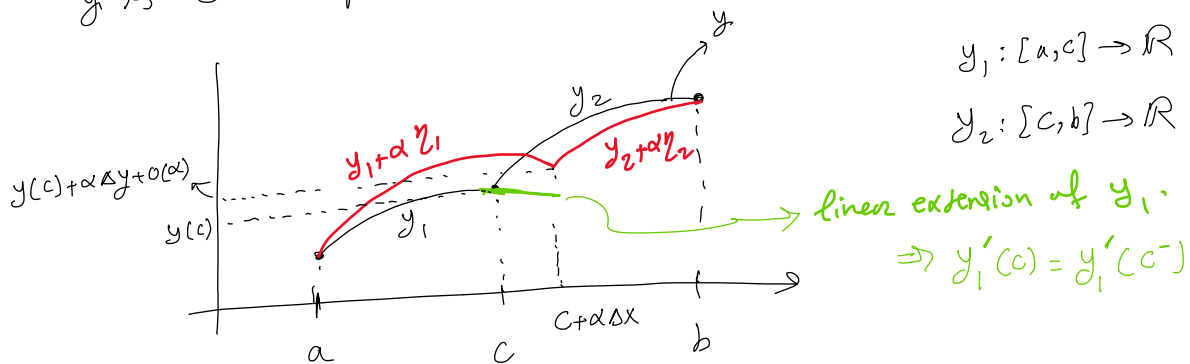
$E \geq 0$ reads that $H^*(z) = H^*(x, y, z, p)$ is
 maximized at $z = y'(x)$!

Note: The Weierstrass condition for strong minimum is so
 powerful that implies the Weierstrass-Erdmann corner
 conditions and the Legendre's condition as corollaries.

Exercise: show this! see Exercise 3.3.

→ let's try to argue the Weierstrass-Erdmann corner conditions:

y is C^1 except at one point $c \in [a, b]$.



→ $y_1 + \alpha \eta_1$ is the perturbed curve on $[a, c + \alpha \Delta x]$ where y_1 is extended linearly beyond c (green line) and

$$\underline{\eta_1(a) = 0}$$

→ $y_2 + \alpha \eta_2$ is defined similarly and $\underline{\eta_2(b) = 0}$

→ To get a piecewise C^1 curve:

$$\textcircled{\text{I}} \quad y_1(c + \alpha \Delta x) + \alpha \eta_1(c + \alpha \Delta x) = y_2(c + \alpha \Delta x) + \alpha \eta_2(c + \alpha \Delta x).$$

$$\triangleq y(c) + \alpha \underline{\Delta y} + o(\alpha).$$

where $\underline{\Delta y}$ is defined, and is independent of Δx .

* Note: we have perturbations $\underline{\eta_1}, \underline{\eta_2}, \underline{\Delta x}, \underline{\Delta y}$ and they are all conditioned on these equalities to be
 ... we will simplify these later!

admissible : ...

* Note: the entire curve is parameterized by the same α ,
and we denote it by $y(\cdot, \alpha): [a, b] \rightarrow \mathbb{R}$
which is C^1 except at $c + \Delta x$.

→ Let's consider the first variations of J with respect to $y(\cdot, \alpha)$:

$$J(y(\cdot, \alpha)) = \underbrace{\int_a^{c+\Delta x} L(x, y_1, y_1') dx}_{\cong J_1(y_1 + \alpha \eta_1)} + \underbrace{\int_{c+\Delta x}^b L(x, y_2, y_2') dx}_{\cong J_2(y_2 + \alpha \eta_2)}$$

Use Leibniz rule to get:

$$\begin{aligned} \delta J_1|_{y_1}(\eta_1) &= \left. \frac{d}{d\alpha} \right|_{\alpha=0} J_1(y_1 + \alpha \eta_1) \\ &= \int_a^c \left(L_y|_{y_1} \eta_1 + L_{y_1'}|_{y_1} \eta_1' \right) dx + L(c, y(c), y'(c)) \cdot \Delta x \end{aligned}$$

Integration
by parts
and

$$\begin{cases} \eta_1(c) \neq 0 \\ y_1'(c) = y_1'(c^-) \end{cases}$$

$$\begin{aligned} &= \int_a^b \left(L_y(x, y_1, y_1') - \frac{d}{dx} L_{y_1'}(x, y_1, y_1') \right) \cdot \eta_1 dx \\ &\quad + L_2(c, y_1(c), y_1'(c^-)) \eta_1(c) + L(c, y_1(c), y_1'(c^-)) \cdot \Delta x \end{aligned}$$

Similarly,

$$\delta J_2|_{y_2}(\eta_2) = \int_c^b \left(L_y(x, y_2, y_2') - \frac{d}{dx} L_{y_2'}(x, y_2, y_2') \right) \cdot \eta_2 dx$$

$$- L_2(c, y(c), y'(c^+)) \eta_2(c) - L(c, y(c), y'(c^+)) \Delta x.$$

So

$$0 = \frac{d}{d\alpha} \Big|_{\alpha=0} \mathcal{J}(y(\cdot, \alpha)) = \frac{d}{d\alpha} \Big|_{\alpha=0} \left(\mathcal{J}_1(y_1 + \alpha \eta_1) + \mathcal{J}_2(y_2 + \alpha \eta_2) \right)$$

$$= \delta \mathcal{J}_1|_{y_1}(\eta_1) + \delta \mathcal{J}_2|_{y_2}(\eta_2),$$

Ⓐ

admissible perturbations $(\eta_1, \eta_2, \Delta x, \Delta y)$.

→ choose the subset of perturbations with $\Delta x = \Delta y = \eta_1(c) = \eta_2(c) = 0$,

then we must have E-L eq for y_1 on $[a, c]$ and y_2 on $[c, b]$.

→ but, now that y_1 and y_2 should satisfy E-L. eq. for all perturbations, condition Ⓐ reads:

$$\begin{aligned} \text{Ⓐ} \quad & L_2(c, y(c), y'(c^-)) \eta_1(c) - L_2(c, y(c), y'(c^+)) \eta_2(c) \\ & + L(c, y(c), y'(c^-)) \Delta x - L(c, y(c), y'(c^+)) \Delta x = 0 \\ & \forall \text{ admissible } (\eta_1, \eta_2, \Delta x, \Delta y). \end{aligned}$$

→ now, we need to characterize the admissible set better from Ⓐ:
look at first variations of Ⓐ w.r.t. α at 0:

$$\frac{y_1'(c) = y_1'(c^-)}{\hookrightarrow} \boxed{y_1'(c^-) \Delta x + \eta_1(c) = y_1'(c^+) \Delta x + \eta_2(c) = \Delta y}$$

this is necessary for perturbations to be admissible,
but is enough to eliminate η_1, η_2 from (III) to get:

$$- L_z(x, y(x), y'(x)) \Big|_{c^-}^{c^+} \cdot \Delta y \\ + \left(L_z(x, y(x), y'(x)) \cdot y'(x) - L(x, y(x), y'(x)) \right) \Big|_{c^-}^{c^+} \cdot \Delta x = 0$$

$\forall$ admissible $\Delta x, \Delta y$ which are independent and totally arbitrary!

$$\Rightarrow L_z \Big|_{c^-}^{c^+} = 0 \text{ and } (L_z \cdot y' - L) \Big|_{c^-}^{c^+} = 0 \quad \square$$